# Adaptive discontinuous Galerkin multiscale methods for elliptic problems:

Energy norm a posteriori error estimate

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#### Outline

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Discontinuous Galerkin (DG)

#### Multiscale method

Framework for Multiscale methods

Discretization

Decay of the fine scale solutions

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## **Papers**

 D. Elfverson, E. Georgoulis and A. Målqvist, Adaptive discontinuous Galerkin multiscale method: Energy norm a posteriori error estimate.

## Model problem

### Poisson's equation

Given a polygonal domain  $\Omega \subset \mathbf{R}^d$ . We want to find u such that

$$\begin{split} -\nabla \cdot \alpha \nabla u &= f \text{ i } \Omega, \\ n \cdot \nabla u &= 0 \text{ on } \partial \Omega, \end{split}$$

where  $\alpha$  is bounded  $0 < \alpha_0 \le \alpha(x) \le \alpha^0$ ,  $f \in L^2(\Omega)$  and  $\int_{\Omega} f \, dx = 0$ .  $L^2$  has the inner product  $(u, v) = \int_{\Omega} uv \, dx$  and norm  $\|u\|^2 = \int_{\Omega} u^2 \, dx$ .

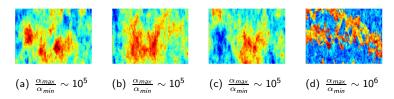


Figure: Permeabilities  $\alpha$  projection in log scale and taken from the Society of Petroleum Engineer http://www.spe.org/

# Weak formulation using a discontinuous Galerkin discretization (SIPG)

- Let  $\Omega$  be subdivided into the partition  $\mathcal{K} = \{\mathcal{K}\}$  and  $\Gamma^I$  be the union of all interior edges.
- Let also  $\mathcal{V}^h$  be the space of all discontinuous piecewise linear polynomials.

The bilinear form and right hand side are defined as:

$$a(v, w) = \sum_{K \in \mathcal{K}} (\alpha \nabla v, \nabla w)_K - \sum_{e \in \Gamma} (\mathbf{n} \cdot \{\alpha \nabla v\}, [w])_e$$
$$- \sum_{e \in \Gamma'} (\mathbf{n} \cdot \{\alpha \nabla w\}, [v])_{\partial K} + \sum_{e \in \Gamma'} \frac{\sigma_e}{h_e} ([v], [w])_e.$$
$$I(v) = (f, v).$$

The discontinuous Galerkin method reads: find  $u_h \in \mathcal{V}^h$  s.t

$$a(u_h, v) = l(v)$$
 for all  $v \in \mathcal{V}^h$ 

## Why do we need to resolve the coefficients?

### Example with periodic coefficient

Consider Possion's equation with period coefficient  $\alpha=\alpha(x/\epsilon)$ . For the finite element method, we have

$$||\sqrt{\alpha}\nabla(u-u_h)||_{L^2(\Omega)}\leq C\frac{H}{\epsilon}||f||_{L^2(\Omega)}$$

- Need  $H \ll \epsilon$  for reliable results.
- To computational expensive to solve on a single mesh for many applications e.g flow in porous media and in composite materials.
- Want eliminate the  $\epsilon$  dependence by using a multiscale method (Målqvist-Peterseim).

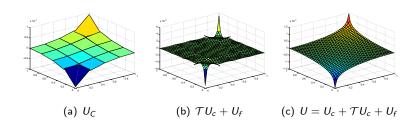
#### Framework for Multiscale methods

The problem is split into one coarse and fine scale contribution  $\mathcal{V}^h = \mathcal{V}_c \oplus \mathcal{V}_f^h$ .

- Let subdivide  $\Omega$  into a coarse mesh  $\mathcal{K}_c = \{\mathcal{K}_c\}$ .
- $\mathcal{V}_c = span\{\phi_i\} = \mathcal{I}_c \mathcal{V}^h$  and  $\mathcal{V}_f^h = \{v \in \mathcal{V}^h : \mathcal{P}_c v = 0\}$ , where  $\mathcal{P}_c : \mathcal{V}^h \to \mathcal{V}_c$  is the  $L^2$  projection onto the coarse mesh.
- Define the map  $\mathcal{T}: \mathcal{V}_c \to \mathcal{V}_f^h$  as  $a(\mathcal{T}v_c, v_f) = -a(v_c, v_f)$ .

Split  $U = U_c + \mathcal{T}U_c + U_f$  and  $v = v_c + v_f$  where  $u_c \in \mathcal{V}_c$ ,  $v_f \in \mathcal{V}_f^h$ .

$$a(U_c + \mathcal{T}U_c + U_f, v_c + v_f) = I(v_c + v_f)$$
 for all  $v_c \in \mathcal{V}_c$  and  $v_f \in \mathcal{V}_f^h$ 



#### Fine scale

Let  $v_c = 0$  to get the fine scale equations

$$a(\mathcal{T}U_c + U_f, v_f) = I(v_f) - a(U_c, v_f),$$

split into two equations

$$\begin{split} &a(\textit{U}_f,\textit{v}_f) = \textit{I}(\textit{v}_f) \quad \forall \textit{v}_f \in \mathcal{V}_f^h, \\ &a(\mathcal{T}\textit{U}_c,\textit{v}_f) = -a(\textit{U}_c,\textit{v}_f) \quad \forall \textit{v}_f \in \mathcal{V}_f^h. \end{split}$$

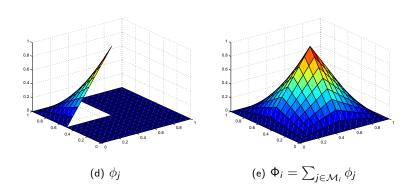
#### Coarse scale

Let  $v_f = 0$  on the coarse scale

$$a(U_c + \mathcal{T}U_c, v_c) = I(v_c) - a(U_f, v_c) \quad \forall v_c \in \mathcal{V}_c$$

## Multiscale method discretization

- ullet  $ilde{\mathcal{T}}$  is he restriction of  $\mathcal{T}$  to a patch  $\omega\subset\Omega$
- $\tilde{U}_f = \sum_{i \in \mathcal{N}} \tilde{U}_{f,i}$  where  $\mathcal{N}$  is the number of nodes.
- $\mathcal{M}_i$  be all j s.t  $\phi_j = 1$  in node i.
- Let also  $\Phi_i = \sum_{j \in \mathcal{M}_i} \phi_j$



### Fine scale equations

For all  $i \in \mathcal{N}$ : find  $\tilde{\mathcal{T}}\phi_j \in \mathcal{V}_f^h(\omega_i^L)$  and  $U_{f,i} \in \mathcal{V}_f^h(\omega_i^L)$  for  $j \in \mathcal{M}_i$  s.t

$$\begin{aligned} & a(\tilde{\mathcal{T}}\phi_j, v_f) = -a(\phi_j, v_f), \quad \forall v_f \in \mathcal{V}_f^h(\omega_i^L), \\ & a(\tilde{U}_{f,i}, v_f) = I(\Phi_i v_f), \quad \forall v_f \in \mathcal{V}_f^h(\omega_i^L). \end{aligned}$$

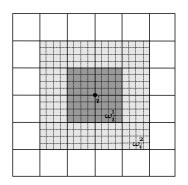
### Coarse scale equation

Find  $U_c \in \mathcal{V}_c$  s.t

$$\textit{a}(\textit{U}_{\textit{c}} + \tilde{\mathcal{T}}\textit{U}_{\textit{c}}, \textit{v}_{\textit{c}}) = \textit{I}(\textit{v}_{\textit{c}}) - (\tilde{\textit{U}}_{\textit{f}}, \textit{v}_{\textit{c}}), \quad \forall \textit{v}_{\textit{c}} \in \mathcal{V}_{\textit{c}}.$$

# Approximation of $TU_c$ and $U_f$

Because of the local behavior in  $\mathcal{V}_f^h$ ,  $\mathcal{T}U_c$  and  $U_f$  can be solved on a subset  $\omega_i^L$  instead of the whole domain  $\Omega$  i.e.  $V_f^h(\omega_i^L) \subset V_f^h$  (Målqvist-Peterseim).



## Decay in $V_f$

#### Problem setting

- Let the computational domain be  $\omega_i^L$  for  $L=1,2,\ldots,N$  where  $\omega_i^L\subseteq\Omega$ .
- Let also  $\Phi_i = \sum_{j \in \mathcal{M}_i} \phi_j$
- ullet The problem reads: find  $ilde{\mathcal{T}}\Phi_i\in\mathcal{V}^h(\omega_i^L)$

$$a(\tilde{\mathcal{T}}\Phi_i, v) = -a(\Phi_i, v), \quad \forall v \in \mathcal{V}^h(\omega_i^L).$$

- The reference solution  $\mathcal{T}\Phi_i$  is the solution computed on  $\omega_i^N=\Omega$ .
- ullet Coarse mesh is 8 imes 8 element and reference grid is 64 imes 64 elements.



Figure: Permeabilities  $\alpha$ .

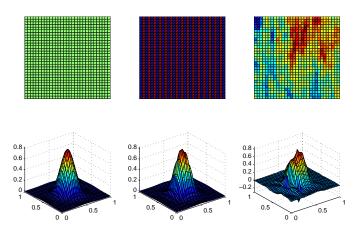


Figure: Example of  $\tilde{T}\Phi_i + \Phi_i$  computed on 2 layer patches.

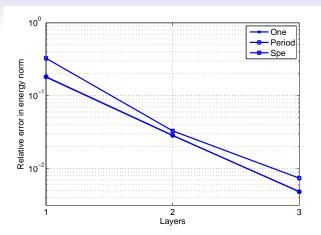


Figure: The error in relative error in broken energy norm with respect to the path size.

## Convergence

### Problem setting

- Consider the model problem (Poisson's equation)
- Keeping the refinement level constant and increasing the patch sizes L = 1, ..., N for all local problems.
- The coarse grid is  $8 \times 8$  coarse elements.
- The reference solution  $U_{ref}$  is the DG solution computed on 64  $\times$  64 elements.
- ullet The right hand side is -1 in the lower left corner and 1 in the upper right.



Figure: Permeabilities  $\alpha$ .

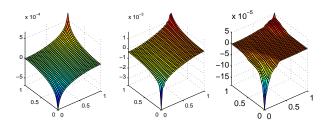


Figure: The reference solution to the model problem using the permeabilities One, Period and SPE

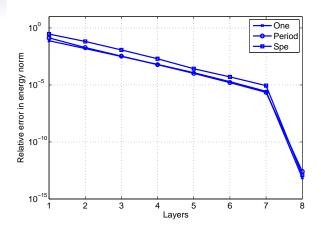


Figure: The relative error in broken energy norm with respect to the patch sizes.

## A posteriori error estimate

- Let  $\mathscr{E} = u U$  where  $U = U_c + \tilde{\mathcal{T}}U_c + U_f$ .
- $U_i = \sum_{j \in \mathcal{M}_i} U_{c,j} (\phi_j + \tilde{\mathcal{T}} \phi_j) + U_{f,i}$ .
- $\lesssim \Leftrightarrow \leq c$  when the constant c is independent of H, h and L.

## Theorem (A posteriori error estimate)

The error & satisfies the estimate

$$\big(\sum_{K \in \mathcal{K}} ||\sqrt{\alpha} \nabla \mathcal{E}||_{L^2(K)}^2\big)^{1/2} \lesssim \sum_{K \in \mathcal{K}_c} \rho_{h,K}^2 + \sum_{i \in \mathcal{N}} \rho_{L,\omega_i^t}^2,$$

where

$$\begin{split} \rho_{L,\omega_i^L}^2 &= \sum_{e \in \Gamma^B(\omega_i^L) \backslash \Gamma^B} \rho_{L,\omega_i^L,e}^2, \\ \rho_{L,\omega_i^L,e} &= \frac{H_{\omega_i^L}}{\sqrt{h_e \alpha_0}} ||n \cdot \{\alpha \nabla U_i\}||_{L^2(e)} + \frac{H_{\omega_i^L} \sqrt{\sigma_e}}{h_e^{3/2}} ||[U_i]||_{L^2(e)}, \end{split}$$

• L is proportional to  $\log(h^{-1})$ .

and

$$\begin{split} \rho_{h,K}^2 &= \sum_{K \in \mathcal{K}_c} \frac{h}{\sqrt{\alpha_0}} ||f + \nabla \cdot \alpha \nabla U||_{L^2(K)} \\ &+ \sum_{e \in \Gamma^I(\mathcal{K}_c)} \left[ \sqrt{\frac{h_K}{\alpha_0}} ||n \cdot [\alpha \nabla U]||_{L^2(e)} + \sqrt{\frac{\sigma}{h_K}} ||[U]||_{L^2(e)} \right] \\ &+ \sum_{e \in \Gamma^B(\mathcal{K}_c) \setminus \Gamma^B} \frac{1}{2} \left[ \sqrt{\frac{h_K}{\alpha_0}} ||n \cdot [\alpha \nabla U]||_{L^2(e)} + \sqrt{\frac{\sigma}{h_K}} ||[U]||_{L^2(e)} \right]. \\ &+ \sum_{e \in \Gamma^B(\mathcal{K}_c) \cap \Gamma^B} \sqrt{\frac{h_K}{\alpha_0}} ||n \cdot \alpha \nabla U||_{L^2(e)}, \end{split}$$

- $ho_{L,\omega^{l}}^{2}$  measure the effect of the truncated patches.
- $\rho_{h,K}^2$  measure the effect of the refinement level.

## Adaptivity

- Consider the model problem
- Using the a posteriori error estimate to construct an adaptive algorithm.
- Start with one refinement and 2 layers patches everywhere.
- Refine 30% of the coarse elements and increase 30% of the patch sizes in each iteration.
- $\bullet$  Coarse mesh is  $32\times32$  element and reference grid is  $256\times256$  elements.
- The right hand side is -1 in the lower left corner and 1 in the upper right.

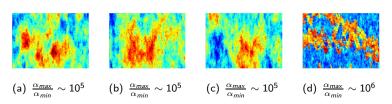


Figure: Permeabilities  $\alpha$  projection in log scale.

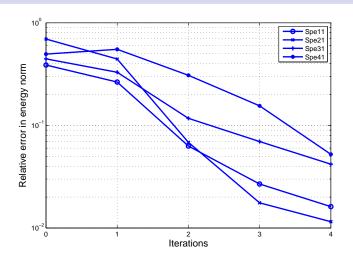
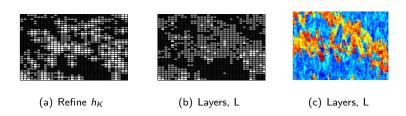


Figure: The relative error in broken energy norm with respect to number of iterations. Iteration 0 corresponds to the standard DG solution and iteration 1 the start values in the adaptive algorithm.

Figure (a) and (b) illustrates where the adaptive algorithm puts most effort

- Figure (a) corresponds to the refinements
- Figure (b) corresponds to the patch sizes.
- Figure (c) is the permeability  $\alpha$ .



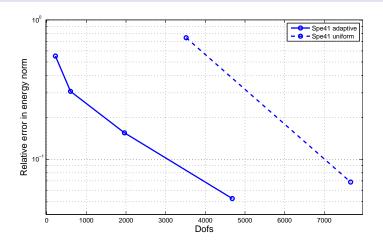


Figure: The relative error in broken energy norm with respect to the mean value of the degrees of freedom for the fine scale problems.

#### Conclusions

### Advantage

- The fine scale problems are perfectly parallelizable.
- The exponential decay in the fine scale solution allows small patches.
- The error estimate and the adaptivity algorithm focus computational effort in critical areas.
- Very high aspect ratio in  $\alpha$  can be solved.
- Possible to construct a conservative flux on the coarse scale.

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#### Future work

- Using a discontinuous Galerkin method with weighted average.
- Solve the local problems iteratively with a appropriate preconditioner. Condition number scale nicely with H,h but not with  $\alpha_{max}/\alpha_{min}$ .
- Convergence of the discontinuous Galerkin multiscale method.
- 3D implementation.

# Questions