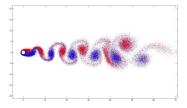
# Parallel Adaptive Fast Multipole Method: application, design, and more...



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PSCP, Uppsala, March 4, 2011



- Background: design of vertical axis wind turbines
- Discretization through vortex formulation
- Fast multipole method...
- ...with adaptivity...
- …and done in parallel

Joint work in part with **Paul Deglaire** and **Anders Goude** at the Division for Electricity and Lightning Research, Uppsala University.

### Background

Pros/cons of VAWTs:

- $+ \ \, {\rm Generator} \ \, {\rm at} \ \, {\rm ground} \ \, {\rm level}$
- + Less gravitational loads
- + No gears
- + Easier maintenance
- + Less noise
  - Fatigue loads
  - Start-up
  - Aerodynamics model

YouTube: Vertical Wind 200kW (March 2010)



## Vortex formulation (the very short version)

In 2D, let the velocity field  $\mathbf{u}(z, t)$  solve the Navier-Stokes equations with BCs (using the complex number z = x + iy for the space coordinate (x, y)). Introduce the *vorticity*  $\omega \equiv \nabla \times \mathbf{u} \cdot \hat{k}$  and consider the two-step formulation:

$$\begin{aligned} \omega_t + \mathbf{u} \cdot \nabla \omega &= 0 \qquad (\text{advection}), \\ \omega_t &= \nu \Delta \omega \quad (\text{diffusion}). \end{aligned}$$

-Hence; how do we obtain **u** from  $\omega$ ?

One can show that  $\mathbf{u} = \mathbf{u}_{\omega} + \nabla \phi$  for some  $\phi$  s.t.  $\Delta \phi = 0$  accounting for the BCs. In turn,

$$\mathbf{u}_{\omega}(z,t) = \int_{\Omega} K(z-z')\omega(z',t) \, dz',$$

where  $K = -i/(2\pi z)$  is the *Green's function* for  $-\Delta$ . If the vorticity is discretized,

$$\omega(z,t)=\sum_{j}\delta(z-z_{j})\Gamma_{j},$$

with  $z_j = z_j(t)$ , then the velocity field is obtained from

$$\mathbf{u}_{\omega}(z,t) = \sum_{j} K(z-z_j) \Gamma_j.$$

To *advect*, evaluate the velocity field in all vorticity points  $z_j$ ,

$$\mathbf{u}_{\omega}(z_j,t) = \sum_{i\neq j} K(z_j-z_i) \Gamma_i,$$

an N-body problem.

To diffuse, just add a normally distributed random number,

$$\mathbf{u}_{\omega}(z_j, t + \Delta t) = \mathbf{u}_{\omega}(z_j, t) + \sqrt{2\nu\Delta t}\mathcal{N}(0, 1).$$

In practice, there are also redistribution-type methods such that  $\Gamma_i$  is made time-dependent.

### Fast multipole method (the *very* short version)

IOLINAL OF COMPUTATIONAL PROJECT 135, 280-292 (1997)

#### A Fast Algorithm for Particle Simulations\*

#### L. Greensard and V. Rokhlin

Department of Computer Science, Yale University, New Haten, Connecticat 06520

#### Received June 10, 1996; nevined February 5, 1997

An algorithm is presented for the rapid evaluation of the potential the potential (or force) at a point is a sum of pairwise and force fields in systems involving large numbers of particles interactions. More specifically, we consider potentials of whose interactions are Coulombic or gravitational in nature. For a the form system of N particles, an amount of work of the order O(N<sup>2</sup>) has less some approximation or truncation method is used. The algo rithm of the present paper requires an amount of work proportional to N to evaluate all interactions to within roundoff error, making it considerably more practical for large-scale problems encountered in mechanics. © 1887 Audemia Press

#### 1 INTRODUCTION

The study of physical systems by means of particle simulations is well established in a number of fields and is becoming increasingly important in others. The most classical example is probably celestial mechanics, but much recent work has been done in formulating and studyine particle models in plasma physics, fluid dynamics, and molecular dynamics [5].

There are two moior classes of simulation methods Dunamical simulations follow the trajectories of N particles over some time interval of interest. Given initial positions (x) and velocities, the trajectory of each particle is gov- of these interactions for all particles. erned by Newton's second law of motion.

$$m_i \frac{d^2 x_i}{dt^2} = -\nabla_i \Phi$$
 for  $i = 1, ..., N$ 

where m is the mass of the ith particle and the force is obtained from the eradient of a potential function  $\Phi$ . When one is interested in an equilibrium configuration of a set of particles rather than their time-dependent properties, an alternative approach is the Monte Carlo method. In this case, the potential function  $\Phi$  has to be evaluated for a large number of configurations in an attempt to determine the potential minimum.

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We restrict our attention in this paper to the case where

#### $\Phi = \Phi_{tor} + (\Phi_{max} + \Phi_{max}),$

where  $\Phi_{max}$  (when present) is a rapidly decaying potential (e.s., Van der Waals), Quetend (when present) is independent of the number of particles, and  $\Phi_{in}$ , the far-field potential, is Coulombic or eravitational. Such models describe classical celestial mechanics and many problems in plasma physics and molecular dynamics. In the vortex method for incompressible fluid flow calculations [4], an the same formal structure (the stream function and the vorticity are related by Poisson's equation).

In a system of N particles, the calculation of  $\Phi_{new}$  requires an amount of work proportional to N, as does the calculation of  $\Phi_{minut}$ . The decay of the Coulombic or gravitational potential, however, is sufficiently slow that all interactions must be accounted for, resulting in CPU time requirements of the order  $O(N^2)$ . In this paper a method is presented for the rapid (order O(N)) evaluation

There have been a number of previous efforts aimed at reducing the computational complexity of the N-body problem. Particle-in-cell methods [5] have received careful study and are used with much success, most notably in plasma physics. Assuming the potential satisfies Poisson's tional domain and the method proceeds by:

(1) interpolating the source density at mesh points, (2) using a "fast Poisson solver" to obtain potential values on the mesh.

(3) computing the force from the potential and interpolatine to the particle positions.

The complexity of these methods is of the order  $O(N + M \log M)$ , where M is the number of mesh points. The number of mesh points is usually chosen to be propertional to the number of particles, but with a small constant NAM I SCI STAT. COMPLET

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#### A FAST ADAPTIVE MULTIPOLE ALGORITHM FOR PARTICLE SIMULATIONS\*

L CARRIERT, L. GREENGARDS, AND V. ROKHLINS

Abstract. This paper describes an algorithm for the rapid evaluation of the potential and force fields in systems involving large numbers of particles whose interactions are described by Coulomb's law. Unlike previously published schemes, the algorithm of this paper has an asymptotic CPU time estimate of O(N), where N is the number of particles in the simulation, and does not depend on the statistics of the distribution for its efficient performance. The numerical examples we present indicate that it should be an alcorithm of choice in many situations of practical interest

Key weeds. N-body problem, plasma physics, molecular dynamics, vortex method, potential theory

AMS(MOS) subject classifications. 65C20, 65D99, 77F05, 82A71, 70F10, 70F15

1. Introduction. The evaluation of Coulombic and gravitational interactions in large-scale ensembles of particles is an integral part of the numerical simulation of a large number of physical processes. Typical examples include celestial mechanics, plasma simulations, the vortex method in fluid dynamics, and the solution of the Laplace equation via potential theory (see [1]-[3], [8], [10]). In such cases, the potential has the form

#### $\Phi = \Phi_{\text{sstanual}} + \Phi_{\text{local}} + \Phi_{\text{far}}$

where  $\Phi_{\text{invest}}$  is a rapidly decaying function of distance (such as the Van der Waals potential in chemical physics),  $\Phi_{external}$  is a function which is independent of the number and relative positions of the particles (such as an external gravitational field) and  $\Phi_{tar}$  is Coulombic or gravitational.

In the numerical evaluation of fields of the form (1), the cost of computing the terms  $\Phi_{n_1, \dots, n_n}$  and  $\Phi_{n_1, \dots, n_n}$  is of the order O(N), where N is the number of particles in the ensemble. Indeed,  $\Phi_{enveral}$  is evaluated separately for each particle, and  $\Phi_{local}$ decays rapidly, involving the interactions of each particle with a small number of nearest neighbors. Unfortunately, evaluation of the term  $\Phi_{tur}$ , if done directly, requires order  $O(N^2)$  operations, since the Coulombic potential decays slowly, and the interactions between each nair of narticles have to be taken into account. In many situations, in order to be of physical interest, the simulation has to involve thousands of particles (or more), making the estimate  $O(N^2)$  excessive in some cases, and prohibitive in others.

Several different approaches have been used to reduce the cost of the Coulombic part of the computation. For a detailed discussion of these algorithms, we refer the reader to [7] and to the original papers [1], [2], [8], [10]. Here, we just observe that each of the algorithms [1], [2], [7], [8], [10] imposes strong requirements on the statistics of the charge distribution. In particular, the methods of [1], [7], and [8] require that the distribution be reasonably uniform in a square-shaped region of interest, the algorithm of [10] assumes that the charges are located on a curve in R<sup>2</sup>, and the algorithm of [2] works fairly well for highly clustered distributions, but fails for uniform ones

\* Received by the editors February 2, 1987; accepted for publication October 8, 1987

+ Elf Aquitaine and Department of Computer Science, Yale University, New Haven, Connecticut 06520. 2 Department of Computer Science, Yale University, New Haven, Connectical 06520. The work of this author was supported in part by the Office of Naval Research under grant N00014-82-K-0184.

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## Fast multipole method (cont)

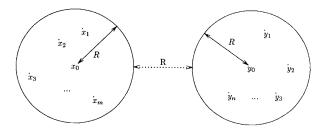


FIG. 1. Well-separated sets in the plane.

Figure: Found at p. 3 of Greengard and Rokhlin: "A Fast Algorithm for Particle Simulations" *J. Comput. Phys.* **73**(2):325–348 (1987).

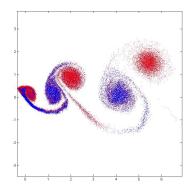
### Fast multipole method (cont)

**Main idea:** all charges/potentials/bodies inside two well-separated sets can interact through an operator of low effective rank. *In a nutshell:* distribute the points in a recursive tree of boxes where each

box has 4 children (2D).

- 1. *Initialize* at the finest level in the tree, expanding each potential in a *multipole series* around the midpoint of the box.
- 2. Go upwards and shift all expansions to parents, yielding a "top expansion" for the whole enclosing box.
- 3. Go downwards and shift-and-convert all expansions into local expansions (eg. polynomials). Also, shift all such expansions to children, yielding a local field in each box.

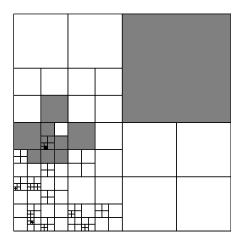
### Illustrations



Test: flat plate. Production run: 3-bladed turbine.

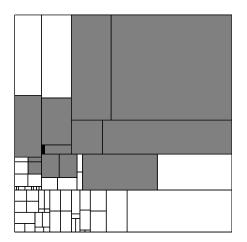
# Adaptivity

Want adaptivity, but quite complicated... The "C" in  $\mathcal{O}(N)$  can be rather large.



### Asymmetric adaptivity

*Idea:* split around the median point instead of around the geometric midpoint. Easier to get the communication localized.



As the mesh looses regularity, it becomes important to keep track of what sets are really well-separated.

#### Criterion

Let the sets  $S_1, S_2 \subset \mathbf{R}^D$  be contained inside two disjoint spheres such that  $||S_1 - x_0|| \le r_1$  and  $||S_2 - y_0|| \le r_2$ . Given  $\theta \in (0, 1)$ , if  $d \equiv ||x_0 - y_0||$ ,  $R \equiv \max\{r_1, r_2\}$ , and  $r \equiv \min\{r_1, r_2\}$ , then the two sets are *well-separated* whenever  $R + \theta r \le \theta d$ .

In other words: any of the two sets may be expanded by a factor of  $1/\theta$  and arbitrarily rotated about its center point without touching the other set.

# Asymmetric adaptivity (cont)

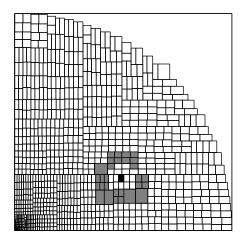


Figure: Typical shift-and-convert interaction list ( $\theta = 1/2$ ).

S. Engblom (UPMARC/TDB)

Parallel Adaptive FMM

PSCP, March 4, 2011 14 / 20

#### Accurate?

*Theory:* the relative error for the *p*th order adaptive fast multipole method under the  $\theta$ -criterion is bounded by a constant  $\times \theta^{p+1}/(1-\theta)^2$ .

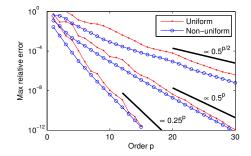


Figure: Errors for two different distribution of points and three distinct  $\theta$ s.

Efficient?

Theory:  $\mathcal{O}\left(\theta^{-2}\log^{-2}\theta\cdot N\log^2 \text{TOL}\right)$ . ( $\Longrightarrow \theta_{\text{opt}} = \exp(-1) \approx 0.368...$ )

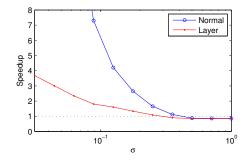
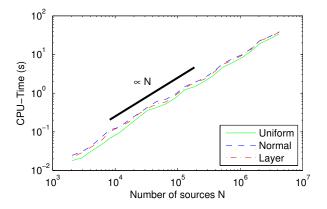


Figure: Adaptive vs. uniform FMM. Two different distribution of points.

"Normal" :=  $\mathcal{N}(0, \sigma)$ , but rejected to fit within the positive unit square. "Layer" := the x-coordinate is U[0, 1] instead.

Parallel Adaptive FMM

### Robustness: does it scale?

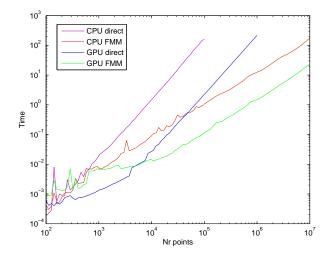


# Implementation at the GPU (ongoing...)

- Started with off-loading the perfectly parallel direct evaluation at the finest level.
- Next came the (not so expensive) local shift operation....
- …and finally the expensive shift-and-convert operation.

What remains serial is currently the initial sorting which is the median-of-three algorithm used in eg. *quicksort. It now takes almost 60–80% of the running time!* 

#### GPU Performance GeForce 480 (NB @ 700MHz!)



### Conclusions

- Academic sw-project: hardly no planning or overall strategy, just a concrete problem to be solved.
- Nice aspect: the steady and controlled growth of performance and complexity:
  - 1. Stand-alone recursive C99 implementation.
  - 2. Matlab-interface to a *direct N*-body evaluation.
  - 3. First working copy; later heavily optimized (eg. BLAS L3).
  - 4. Added adaptivity took the time to investigate a novel approach.
  - 5. Parallel GPU-implementation CUDA/C++ (still ongoing).
  - 6. 3D...?
- Increasingly sophisticated regression tests.
- In such an academic environment, clarity wrt to goals is very important.