



# High-order convergent and accurate electromagnetic solvers on Lipschitz domains

Johan Helsing, Anders Karlsson, and Karl-Mikael Perfekt  
Lund University

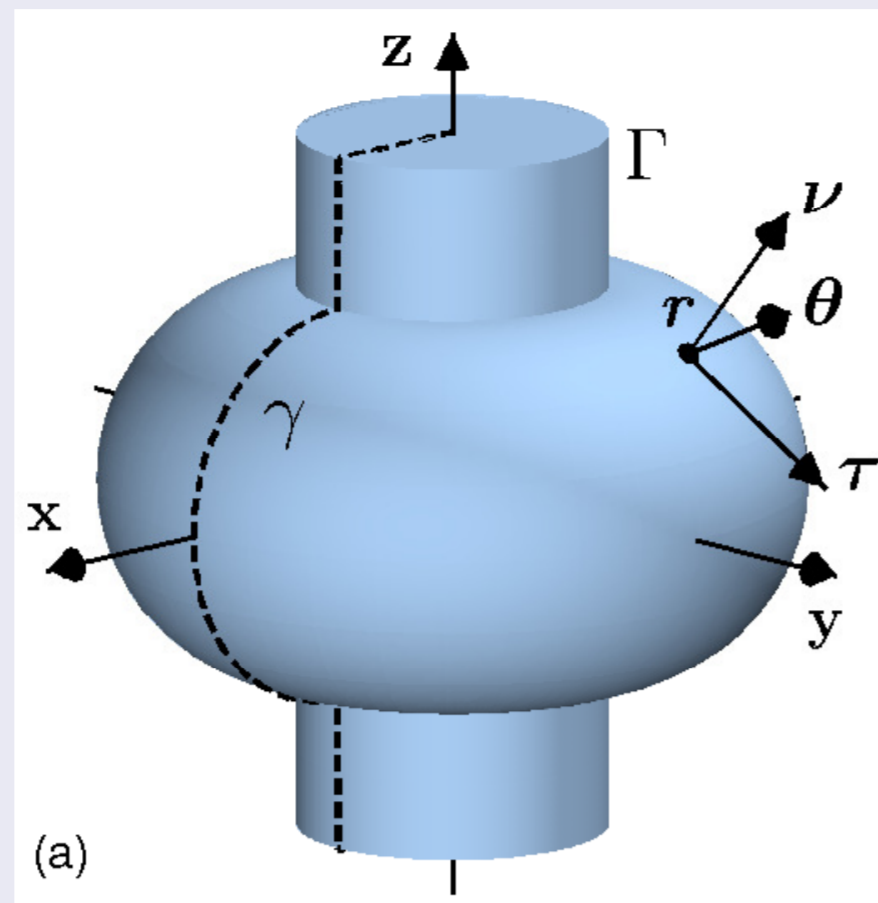
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## Problem, challenges, and methods

**Problem:** Solve the time harmonic Maxwell equations and find eigenwavenumbers  $k$  and eigenfields  $E$  of perfectly conducting axially symmetric cavities and for dielectric objects in vacuum.

**Challenges:** Large structures. High accuracy. Finding all eigenmodes. Resolution of singular fields at sharp edges. Normalization  $\|E\|^2 = 1$  of eigenfields.

**Methods:** ChIE extensions of MFIE and EFIE (no surface divergence for surface charge densities), 16th-order explicit kernel-split panel-based Fourier–Nyström discretization, recursively compressed inverse preconditioning (RCIP), volume→surface integral for  $\|E_n\|^2 = 1$ , robust search for eigenwavenumbers.



## Integral equation for eigenproblem

The ChIE extended MFIE system can be written on modal block operator form as

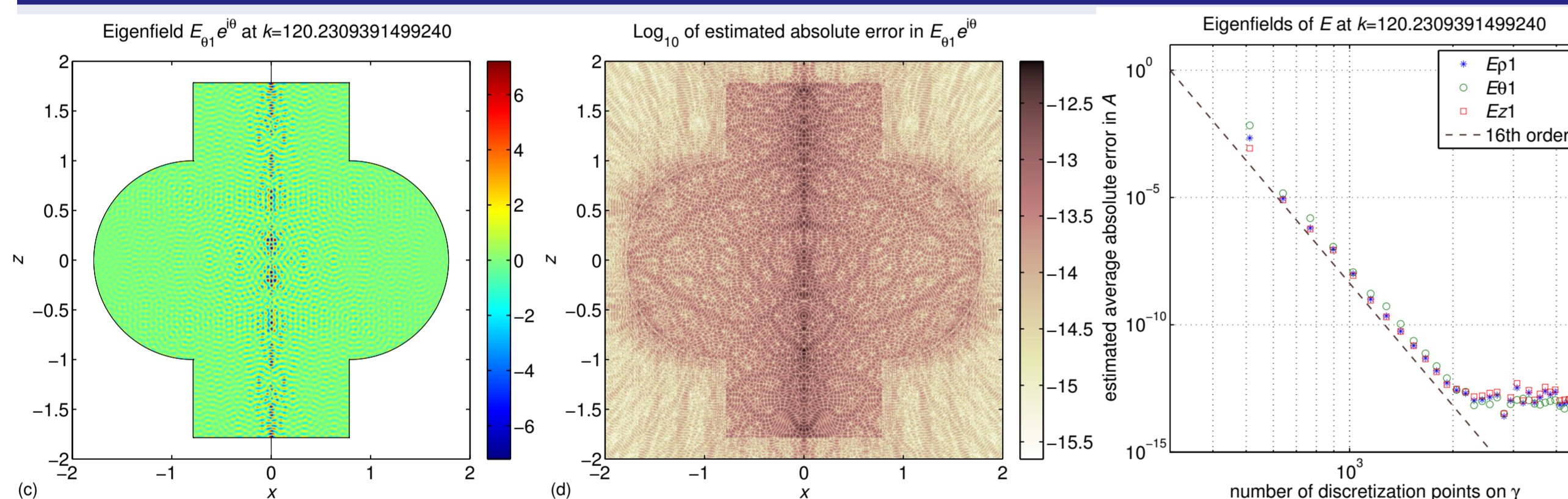
$$\begin{bmatrix} I - 2K_{\nu n} & 2ikS_{5n} & -2kS_{6n} \\ 0 & I + K_{1n} & iK_{2n} \\ 0 & iK_{3n} & I + K_{4n} \end{bmatrix} \begin{bmatrix} Q_{sn}(r) \\ J_{\tau n}(r) \\ J_{\theta n}(r) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix},$$

where  $K_{\nu n}$ ,  $K_{1n}$ ,  $K_{2n}$ ,  $K_{3n}$ ,  $K_{4n}$  are modal double-layer type integral operators and  $S_{5n}$ ,  $S_{6n}$  are single-layer type operators. All operators are weakly singular.

## The RCIP idea

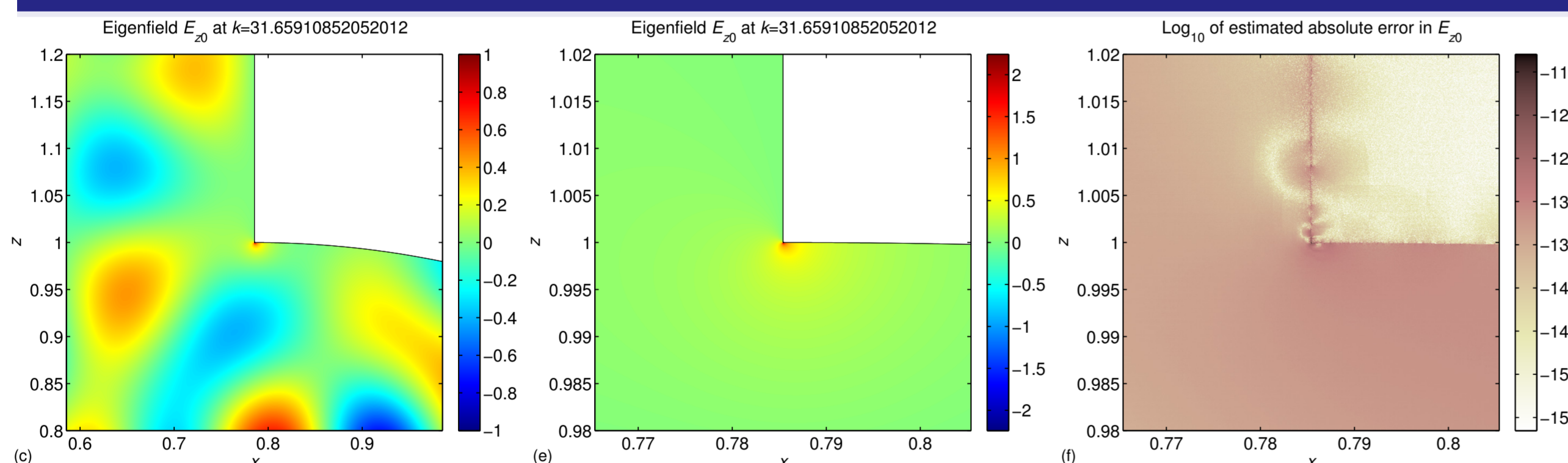
$(I + K)\rho = g$  is difficult to solve on Lipschitz domains. The split  $K = K^* + K^\circ$  and the change of variables  $\rho = (I + K^*)^{-1}\tilde{\rho}$  give the simpler compressed preconditioned equation  $(I + K^\circ P_W^T (I + K^*)^{-1} P)\tilde{\rho} = g$ , where  $P$  is a prolongation operator. Recursion on nested grids is used for the lossless compression of  $P_W^T (I + K^*)^{-1} P$ .

## One cell elliptic cavity with edges



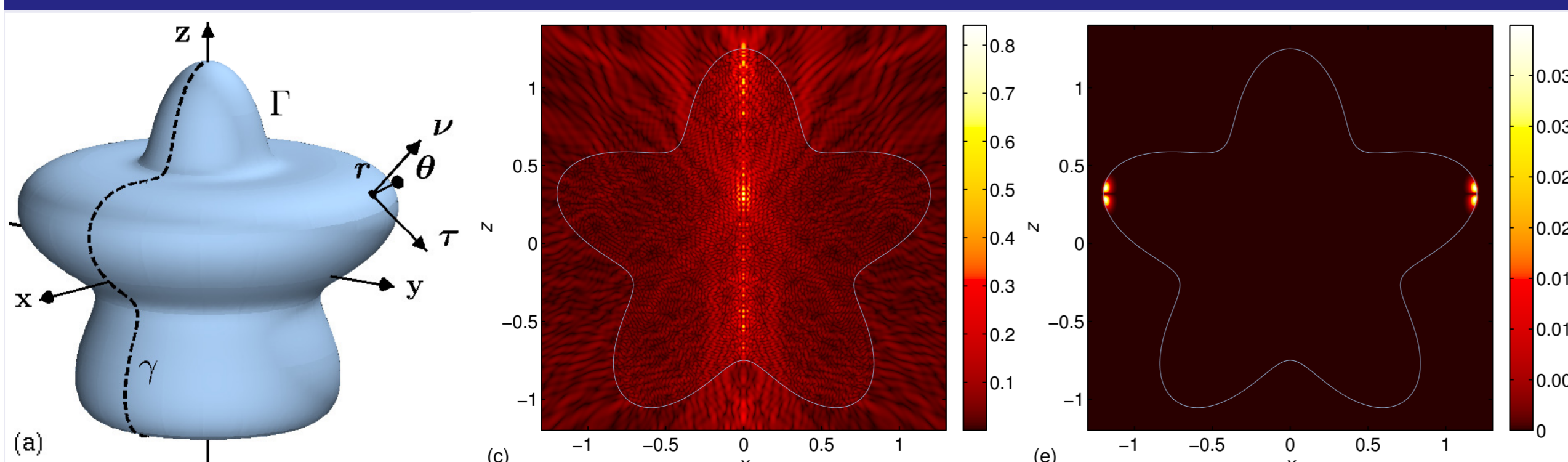
Left: The azimuthal component of normalized electric field for the 9928th  $n = 1$  eigenmode with eigenwavenumber  $k = 120.2309391499240$  evaluated at  $4.9 \cdot 10^5$  points on a cartesian grid. Object diameter  $\approx 75\lambda$ . Center:  $\log_{10}$  of estimated absolute pointwise error. Right: stable 16th-order convergence with mesh refinement.

## Magnification of edge singularity



Left and center: zoom of  $z$ -component of normalized electric field for the 662nd  $n = 0$  eigenmode with eigenwavenumber  $k = 31.65910852052012$ . The field diverges in the reentrant corners. Right:  $\log_{10}$  of estimated absolute pointwise error.

## Dielectric object in vacuum



Left: Object. Center: Azimuthal component of electric field for the  $n = 1$  eigenmode with eigenwavenumber  $k = 110.041232211051 - 0.404177078290i$ . The refractive index is  $m = 1.5$ . The object diameter is around 46 vacuum wavelengths. Right: Whispering gallery  $n = 450$  mode with eigenwavenumber  $k = 258.059066513439$ ,  $m = 1.5$ , and object diameter around 108 vacuum wavelengths. The absolute pointwise error is less than  $10^{-11}$  in both field plots.

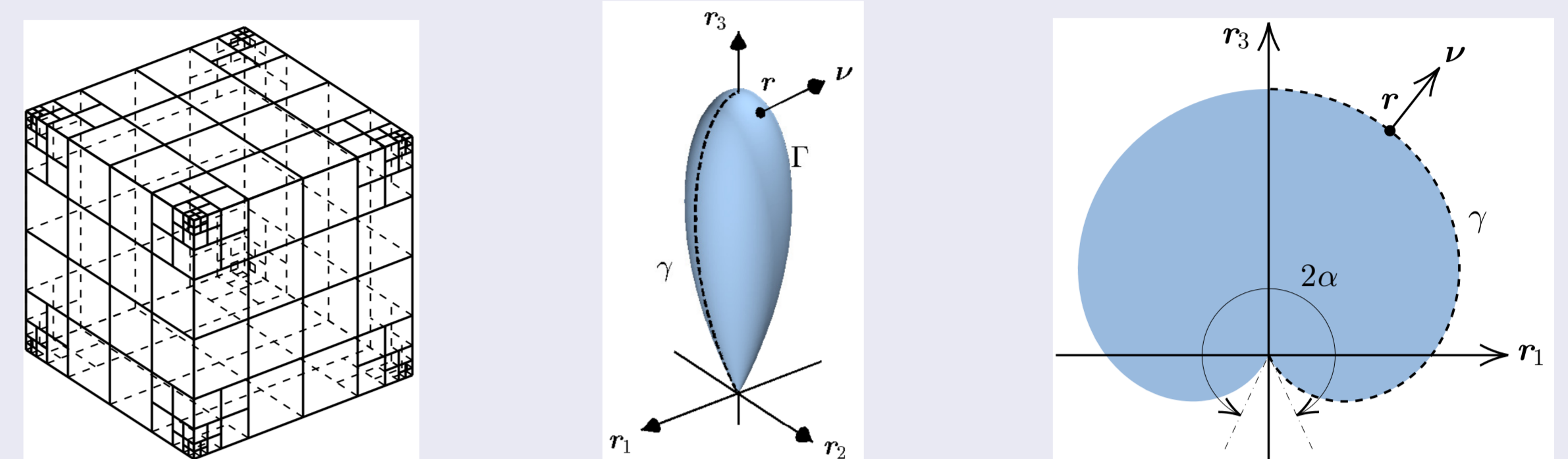
## Problem, challenges and methods

**Problem:** Solve the electrostatic transmission problem for an inclusion with Lipschitz surface  $\Gamma$  and permittivity  $\epsilon$  embedded in a background medium with unit permittivity. Then compute the polarizability tensor  $\omega_{ij}(\epsilon)$  and its spectral measure.

**Challenges:** High accuracy. Resolution of singular fields at sharp edges and in corners. Integral operator spectra depend on function spaces considered:  $L^2(\Gamma)$  or energy space.

**Methods:** Classic integral equation, 16th- and 32nd-order explicit kernel-split panel-based (Fourier–)Nyström discretization, recursively compressed inverse preconditioning (RCIP), fixed-point iteration, Newton's method, homotopy.

## Geometries and meshes



Left: a three times dyadically refined mesh on the surface of a cube. Middle: An axially symmetric surface  $\Gamma$  with a conical point of opening angle  $\alpha = 5\pi/36$ , denoted a snow cone. Right: A cross section of the snow cone interior for  $\alpha = 31\pi/36$ .

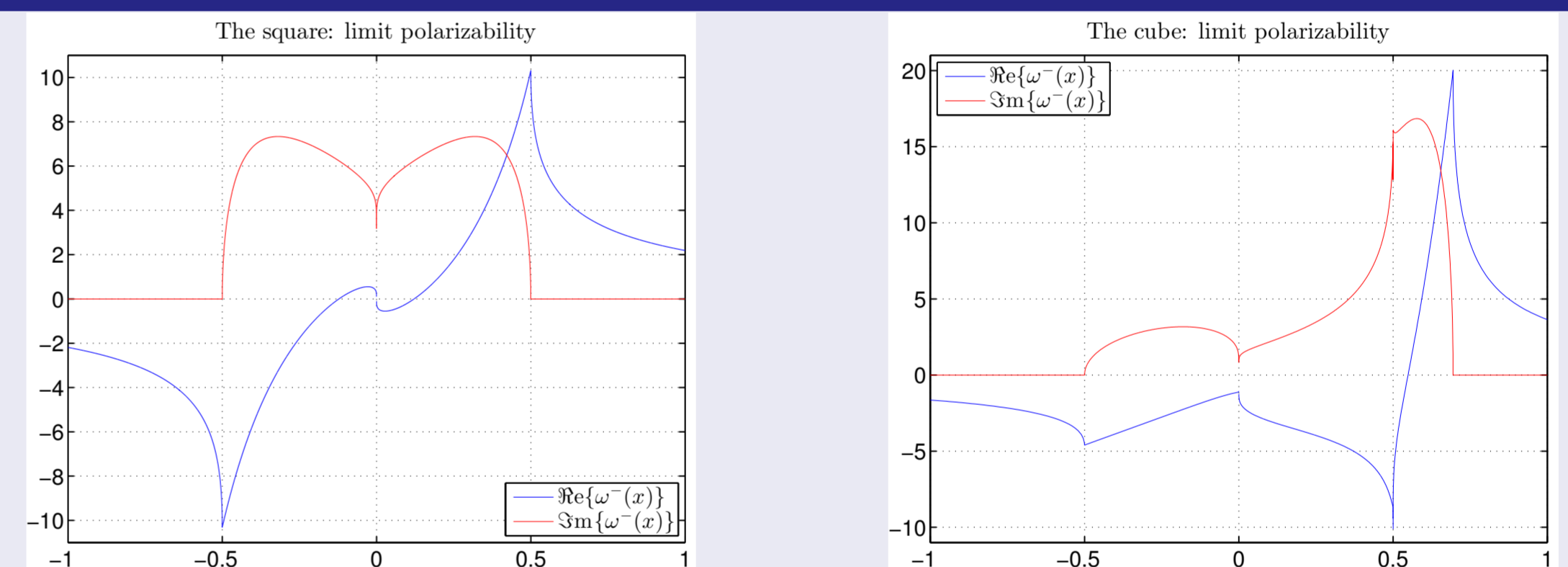
## Integral equation

With standard basis  $e_i$ , the element  $\omega_{ij}$  of the polarizability tensor of an inclusion is

$$\omega_{ij}(z) = \frac{2}{|V|} \int_{\Gamma} (e_i \cdot r) [(K_{\nu} - z)^{-1} (e_j \cdot \nu)](r) d\sigma(r),$$

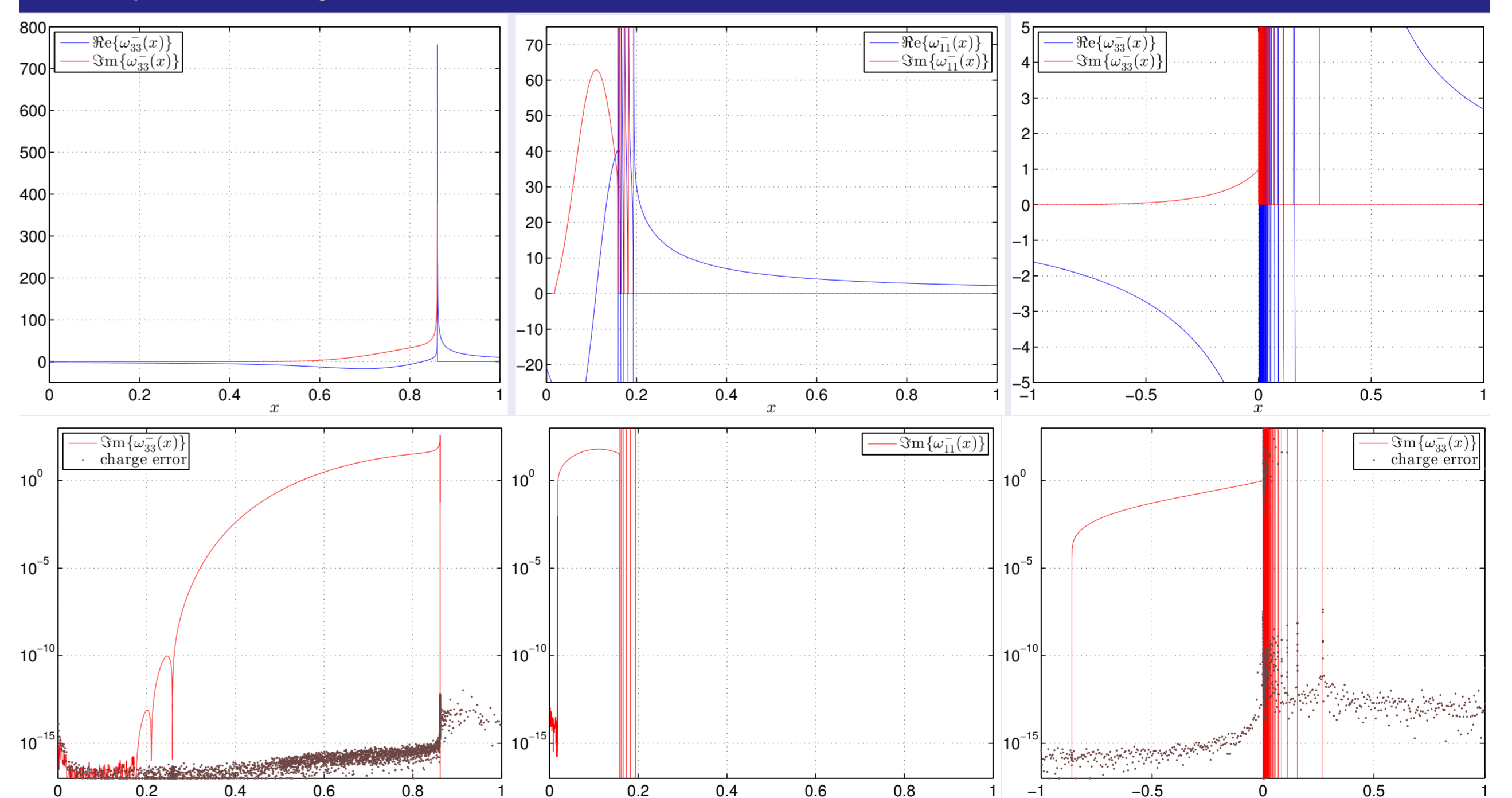
where  $K_{\nu}$  is two times the adjoint of the Neumann–Poincaré operator (double layer potential),  $z \equiv x + iy = (\epsilon + 1)/(\epsilon - 1)$ , and  $|V|$  is the volume of the inclusion.

## Limit polarizability of square and cube



Limit values of  $\omega(z)$  as  $y \rightarrow 0^-$ . The accuracy varies from full machine precision to about five digits.  $\Im m\{\omega^-(x)\}$  of the cube has support for  $x \in (-0.5, 0.694526)$ .

## Limit polarizability of snow cones



Top row: (a,b) limit values of  $\omega_{33}(z)$  and  $\omega_{11}(z)$  as  $y \rightarrow 0^-$  for  $\alpha = 5\pi/36$ ; (c)  $\omega_{33}(z)$  for  $\alpha = 31\pi/36$ . Bottom row: imaginary parts with logarithmic scales on the vertical axes. The spectral measure  $\mu_i(x)$ , associated with  $\omega_{ii}(z)$ , is determined by  $\mu_i'(x) = -\Im m\{\omega_{ii}^-(x)\}/\pi$ . The numerical accuracy in (a,b) is such that  $\int_{-1}^1 d\mu_i(s) = -2$  holds to almost machine precision. Note the infinite number of poles for  $\omega_{33}^-(x)$ ,  $x > 0$ , and  $\alpha = 31\pi/36$ , of which 275 are located and drawn.

## References

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